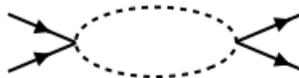


The S-matrix Bootstrap

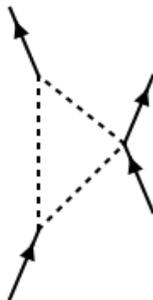
Problem sheet 1

1. Show that the bubble diagram



has a square root cut, i.e. behaves like $\sqrt{s - 4m_\chi^2}$ near $s = 4m_\chi^2$.

2. Consider a theory with $\mathcal{L} = \frac{1}{4}g\phi^2\chi^2$ in $d = 4$. Using your previous answer, give the ϕ - ϕ elastic scattering amplitude to second order in g . Use the partial wave decomposition to verify that the unitarity condition is satisfied up to order g^2 .
3. Consider the elastic scattering of two real scalar particles of mass m^2 in $d = 4$. Suppose the amplitude is real-analytic and saturates unitarity for a finite range of s near threshold. Prove that the amplitude must have a two-sheeted branch cut starting at $s = 4m^2$.
4. Compute the location of the anomalous threshold singularity in the real (s, t) plane for the Landau diagram:



You may suppose that all momenta are Euclidean and planar. Hint: make a simple ansatz for the Euclidean momenta like $p_1 = m(\cos(\phi), \sin(\phi))$, etc.

5. Consider scattering amplitudes in *two* spacetime dimensions, again for real identical scalar particles. Using the rapidity variables θ_i , defined such that $p_i = m(\cosh(\theta_i), \sinh(\theta_i))$, show that for on-shell momenta

$$\delta^{(2)}(p_1 + p_2 - p_3 - p_4) \propto \delta(\mathbf{p}_1 - \mathbf{p}_4)\delta(\mathbf{p}_2 - \mathbf{p}_3) + (2 \leftrightarrow 3) \tag{1}$$

and find the exact proportionality constant. Show that we can therefore write

$$\langle \mathbf{p}_3, \mathbf{p}_4 | S | \mathbf{p}_1, \mathbf{p}_2 \rangle = (2\pi)^2 4E_1 E_2 (\delta(\mathbf{p}_1 - \mathbf{p}_4)\delta(\mathbf{p}_2 - \mathbf{p}_3) + (1 \leftrightarrow 2)) S(s) \tag{2}$$

with $s = (p_1 + p_2)^2$ the only remaining Mandelstam invariant. Show that crossing symmetry implies that $S(4m^2 - s) = S(s)$. What is the constraint of unitarity? Give $S(s)$ for a simple contact interaction corresponding to $\mathcal{L}_{int} = \lambda\phi^4/(4!)$.

6. Consider the Virasoro-Shapiro amplitude (in units such that $\alpha' = 1$)

$$T_{VS}(s, t) = \left(\frac{tu}{s} + \frac{su}{t} + \frac{st}{u} \right) \frac{\Gamma(1 - s/4)\Gamma(1 - t/4)\Gamma(1 - u/4)}{\Gamma(1 + s/4)\Gamma(1 + t/4)\Gamma(1 + u/4)} \quad (3)$$

and the partial wave expansion

$$T_{VS}(s, t) = \sum_{\ell=0}^{\infty} f_{\ell}(s) C_{\ell}(\cos(\theta)) \quad (4)$$

where now $\cos(\theta) = 1 + 2t/s$ and the $C_{\ell}(z) = c_{\ell}z^{\ell} + \dots$ are polynomials of degree ℓ ; you can think of these as ten-dimensional versions of the Legendre polynomials. Notice also that the external particles are massless so $s + t + u = 0$.

- (a) Show that the residues of the scattering amplitude at $s = 4n$ are polynomials in $\cos(\theta)$ of degree $2n + 2$. For which values of n does $f_{\ell}(s)$ then have a pole?
- (b) Consider *only* the poles in $f_{\ell}(s)$ at $s = 2\ell - 4$, which form the *leading* Regge trajectory. Use this trajectory together with crossing symmetry to show that

$$T_{VS}(s, t) \sim s^{2+t/2} \quad (5)$$

for large s and fixed t . Check that this is the correct result by explicitly taking the limit of $T_{VS}(s, t)$. Can you match the exact coefficient as well?